

A UNIQUENESS THEOREM OF THE INVERSE PROBLEM FOR A CLASS THE STURM– LIOUVILLE PROBLEM

Kh. R. Mamedov, U. Demirbilek (Mersin, TURKEY)

hanlar@mersin.edu.tr, ulviyedemirbilek@gmail.com

In the present paper , we study inverse problem of scattering theory for Sturm-Liouville operator on the half-axis $[0, \infty)$ with spectral paramater in the boundary condition for second order of differantial equation .We define the kernel function and determine the scattering data uniquely .

Keywords: Sturm–Liouville operator, scattering function, scattering data, Gelfand-Levitan-Marchenko main equation..

Introduction

We consider inverse problem of scattering theory for the Sturm–Liouville equation

$$-y'' + q(x)y = \lambda^2 y, \quad (1)$$

on the semi-axis $[0, \infty)$ containing a spectral parameter in the boundary condition

$$y'(0) + (\alpha_0 + \alpha_1 \lambda + \alpha_2 \lambda^2)y(0) = 0. \quad (2)$$

Here λ is a spectral parameter, α_i ($i = 0, 1, 2.$) are real numbers that satisfy certain conditions. $q(x)$ is a real valued function satisfying the condition

$$\int_0^{\infty} (1+x) |q(x)| dx < \infty. \quad (3)$$

It is well known (see [1]) that for all λ from the half-line Eq. (1) has the solution

$$e(x, \lambda) = e^{i\lambda x} + \int_x^{\infty} K(x, t) e^{i\lambda t} dt. \quad (4)$$

The kernel $K(x, t)$ satisfies the inequality

$$K(x, t) \leq \frac{1}{2} \sigma\left(\frac{x+t}{2}\right) \exp \left\{ \sigma_1(x) - \sigma_1\left(\frac{x+t}{2}\right) \right\}. \quad (5)$$

The inverse problem of scattering data without any spectral parameter in boundary condition was solved in [1, 2]. The many spectral properties of the boundary value problems were investigated with different methods by the many authors in [1–7].

In this work we prove the uniqueness of the solution of the inverse problem of scattering theory on the half line for the boundary – value problem (1)–(2) by using defined the scattering data of the problem and its properties.

With the above preliminaries provided, we have the following lemmas and theorems.

Lemma 1. *For all $\lambda \neq 0$, the identity is valid*

$$\frac{2i\lambda w(x, \lambda)}{E(\lambda)} = e(x, -\lambda) - S(\lambda)e(x, \lambda), \quad (6)$$

where

$$S(\lambda) = \frac{E_1(\lambda)}{E(\lambda)}, \quad (7)$$

$$E(\lambda) = e'(0, \lambda) + (\alpha_0 + \alpha_1\lambda + \alpha_2\lambda^2)e(0, \lambda),$$

$$E_1(\lambda) = e'(0, -\lambda) + (\alpha_0 + \alpha_1\lambda + \alpha_2\lambda^2)e(0, -\lambda)$$

and

$$|S(\lambda)| = 1.$$

The scattering function $S(\lambda)$ is meromorphic in half plane $Im\lambda > 0$, with poles at the zeros of the function $E(\lambda)$. Moreover, the function $E(\lambda)$ is analytic in upper half plane. The function $E(\lambda)$ may have only a finite number of zeros in the half plane $Im\lambda > 0$.

We shall obtain the main equation that contributes to construct the potential $q(x)$ in the Eq. (1). To obtain the main equation, we substituting the relation (4) into the relation (6). Thus, the following results are valid:

Theorem 1. *For each fixed $\neq 0$ the kernel $K(x, t)$ satisfies the following equation:*

$$F(x + y) + K(x + y) + \int_x^\infty K(x, t)F(t + y)dt = 0, \quad x < y < \infty, \quad (8)$$

Proof. From [3, Theorem 3.1.] it is clear that the main equation can be constructed.

Thus, we have the following theorem.

Theorem 2. *For each $x \geq 0$, the kernel $(K(x, t))$ to the solution (4) satisfy the main equation (8).*

Uniqueness

Lemma 2. Assume that the function $f_x(t)$ is summable on the half line for $t \geq x$

$$f_x(t) + \int_x^\infty f_x(u)F(u+t)du = 0, \quad (9)$$

and there is a solution for $f_x(t) \equiv 0$, $t \geq x$.

Proof. It can be easily seen from [3, Lemma 4.1], the function $f_x(t)$ be solution of the integral equation, where $K(x, t)$ satisfies the equation (8). Then, the homogeneous equation (9) has only trival solution i.e. $f_x(t) \equiv 0$ for $t \geq x$.

Theorem 3. The scattering data of the boundary value problem (1)–(3) determine uniquely.

Proof. Given scattering function $S(\lambda)$ for $\lambda \neq 0$ and the scattering data can be determined according to Eq. (8). By virtue of the function $F(x)$, the main equation is constructed and it sufficies to find only scattering data of the boundary value problem (1)–(3). Given the scattering data, we can use formulas as follows:

$$F_s(x) = \frac{1}{2\pi} \int_x^\infty [1 - S(\lambda)]e^{i\lambda x} d\lambda,$$

$$F(x) = \sum_{j=1}^n f_j(x) + F_s(x),$$

and

$$p_j(x) = e^{-i\lambda_j x} f_j(x), \quad j = 1, 2, \dots, n.$$

By Lemma 2, the main equation has a unique solution. Futhermore, we find the function $K(x, t)$. It follows from, applying (5) we have

$$q(x) = -2 \frac{d}{dx} K(x, x). \quad (10)$$

Thus, the potential $q(x)$ can uniquely be found from (10).The theorem is proved.

REFERENCES

- [1] Marchenko V. A. Sturm–Liouville Operators and Their Applications. Kiev : Naukova Dumka, 1977 (in Russian).
- [2] Levitan B. M. Inverse Sturm–Liouville Problems. Utrecht : VNU Science Press, 1987.

- [3] *Mamedov Kh. R.* Uniqueness of the solution of the inverse problem of scattering theory for the Sturm–Liouville operator with a spectral parameter in the boundary condition // *Math. Notes.* 2003. Vol. 74, № 1–2. P. 136–140.
- [4] *Mamedov Kh. R.* On The Inverse Problem For Sturm–Liouville Operator With A Nonlinear Spectral Parameter In The Boundary Condition // *J. Korean Math. Soc.* 2009. Vol. 46, № 6. P. 1243–1254.
- [5] *Yurko V. A.* An inverse problem for pencil of differential operator on the half line // *Sbornik Math.* 2000. Vol. 191, № 9–10. P. 1561–1586.
- [6] *Yurko V. A.* Reconstruction of the pencils of differential operators on the half line // *Mat. Zametki.* 2000, Vol. 67, № 2. P. 316–320; *Math. Notes.* 2000. Vol. 67, № 1–2. P. 261–265..
- [7] *Yang Y., Wei G.* Inverse Scattering Problems for Sturm Liouville Operators Spectral Parameter Dependent on Boundary Condition // *Math. Notes.* 2018. Vol. 103, № 1. P. 65–74.